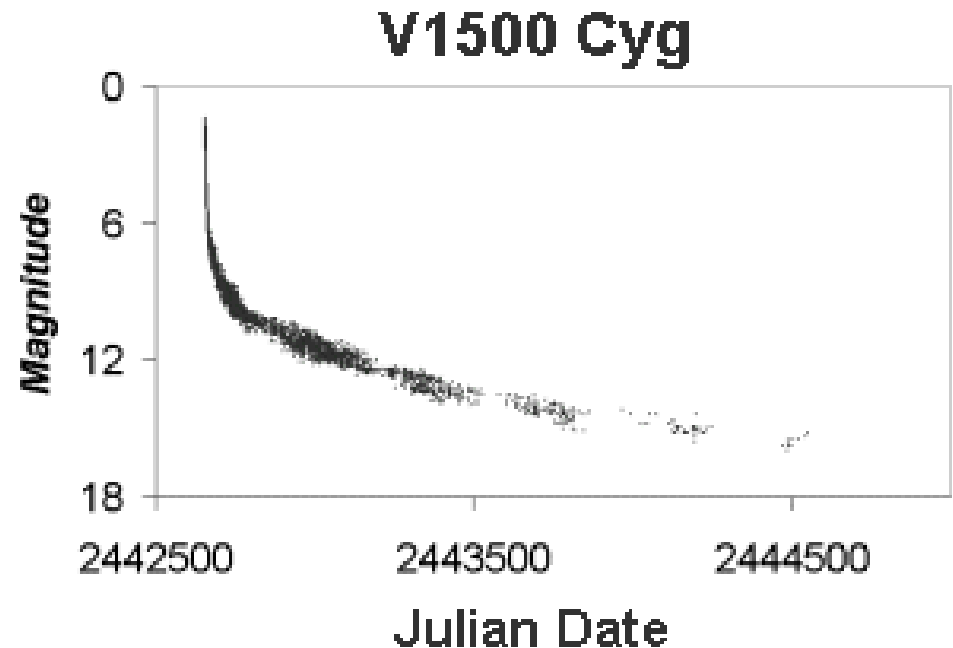


Thermonuclear shell flashes II: on WDs (or: classical novae)

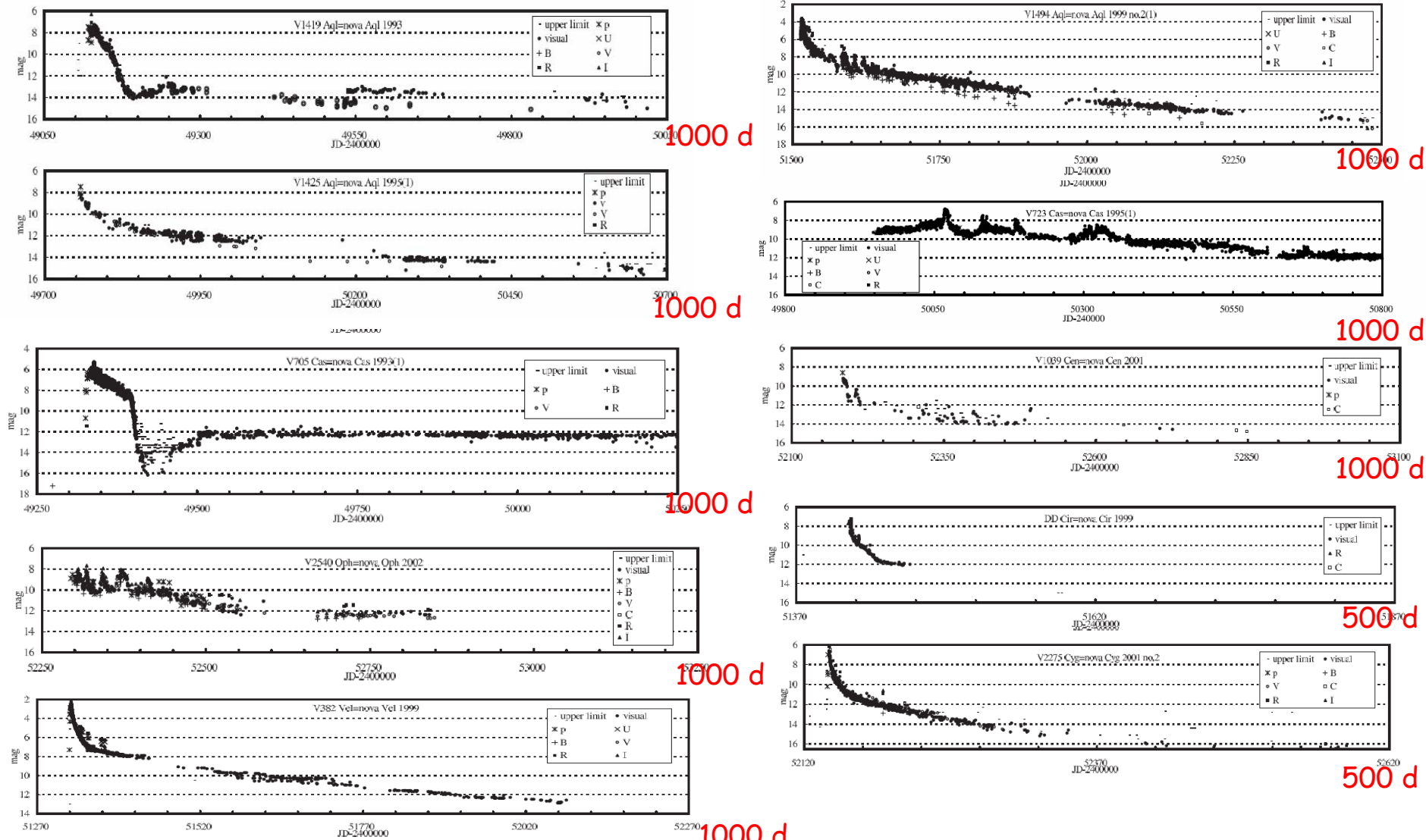
- Observations
- Thermonuclear flash model
- Nova/X-ray burst comparison
- Effects of super-Eddington fluxes
- To grow or not to grow =
to go supernova Ia or not ..

Nova Cygni 1975

- Nova Cygnus 1975 is the 4th brightest nova in modern times (after Nova Persei 1901 [0.2], Nova Aurigae 1918 [-1.8] and Nova Puppis 1942 [0.3]) $V = 2.0$ mag). No brighter nova since then
- Descend of 7 mag in 45 days
- V1500 Cyg turned out to be a magnetic CV with an orbital period of 3.3 hr

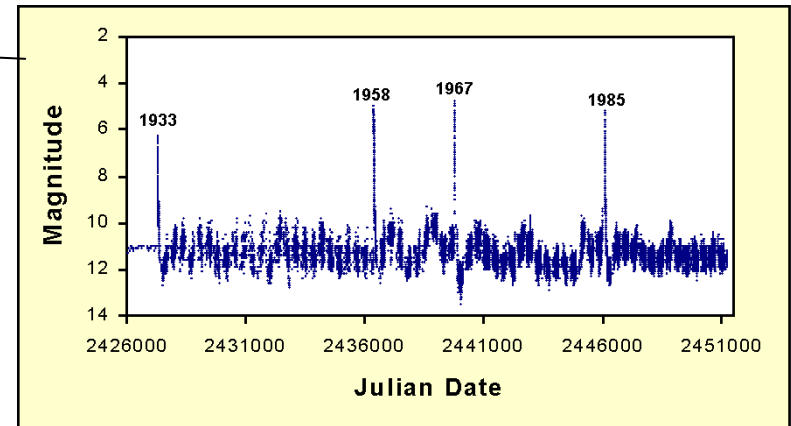
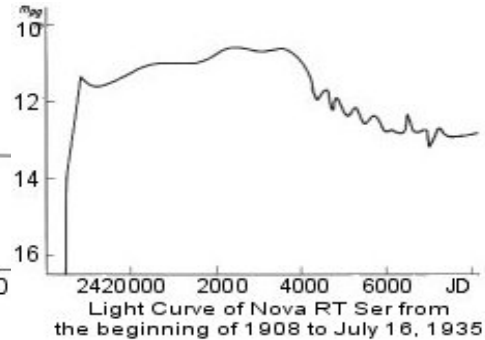
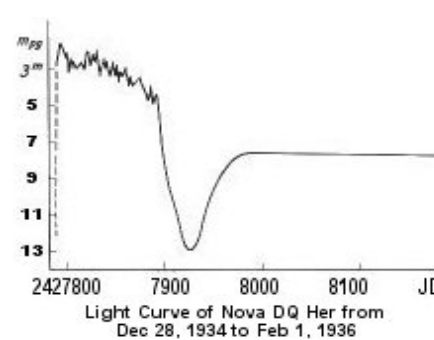
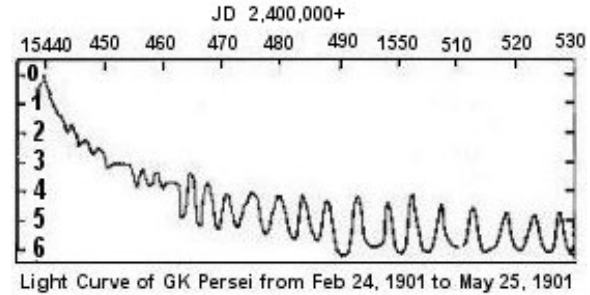


Examples of classical nova light curves

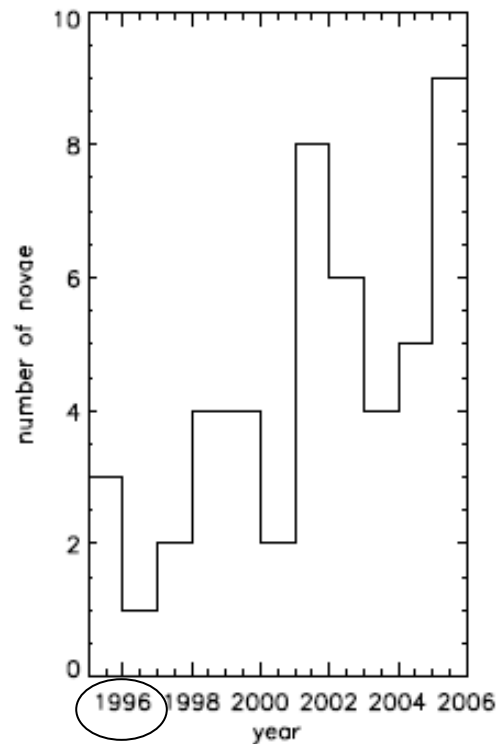
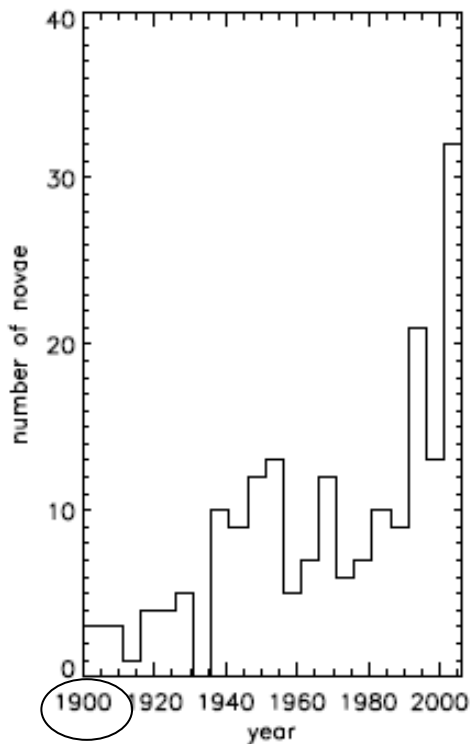


Classical novae characteristics

- Fast ($t_2 < 10$ d) , slow ($10 < t_2 < 25$) and very slow novae ($t_2 > 25$ days) (t_2 time needed to faint 2 magnitudes)
- Peak brightness: $m_V \approx 4-12$ mag
- Peak luminosity: $M_V \approx -6/-11$
- $\Delta V \approx 7-16 \rightarrow L_{\max} / L_{\min}$ up to 10^6
- Rise time few days
- Decay time weeks to decades
- Recurrent novae \rightarrow recurrence observed and thus less than ~ 100 yr
- Light curves and spectra very diverse because of shedding of matter (\rightarrow reprocessing of light in circumstellar environment)
- Most discoveries done by amateurs
- Thus far ~ 250 discovered, 10 recurrent



Nova discovery rates



→ On average 5 per year

V2362	Cyg	N2006	04	02.807	21	11	32.34	+44	48	03.9
V2576	Oph	N2006	04	06.565	17	15	33.00	-29	09	39.9
V1065	Cen	N2007	01	23.354	11	43	10.33	-58	04	04.3
V1280	Sco	N2007	02	04.9	16	57	41.20	-32	20	35.6
V1281	Sco	N2007	02	19.8593	16	56	59.35	-35	21	50.2
V2467	Cyg	N2007	03	15.787	20	28	12.52	+41	48	36.5
V2615	Oph	N2007	03	19.612	17	42	44.00	-23	40	35.1
V5558	Sgr	N2007	04	14.777	18	10	18.27	-18	46	52.1
V390	Nor	N2007	06	15.086	16	32	11.51	-45	09	13.4
V458	Vul	N2007	08	04.64	19	54	24.64	+20	52	51.7
V597	Pup	N2007	08	14.21	08	16	18.01	-34	15	24.1
V598	Pup	N2007	08	08	07	05	42.51	-38	14	39.3
V459	Vul	N2007	12	25.35	19	48	08.84	+21	15	26.8

10.5	I8697	H. Nishimura
10.5	I8700	P. Williams
8.2	I8800	W. Liller
9.4	I8803	Y. Sakurai; Y. Nakamura
9.3	I8810	Y. Nakamura; H. Nishimura
7.4	I8821	A. Tago
10.2	I8824	H. Nishimura
10.3	I8832	Y. Sakurai
9.4	I8850	W. Liller
9.5	I8861	H. Abe
7.0	I8895	A. Pereira
10.3	I8898	Read et al. (x-ray)
8.7	I8907	Kaneda

Go to the [Index](#)

Classical novae

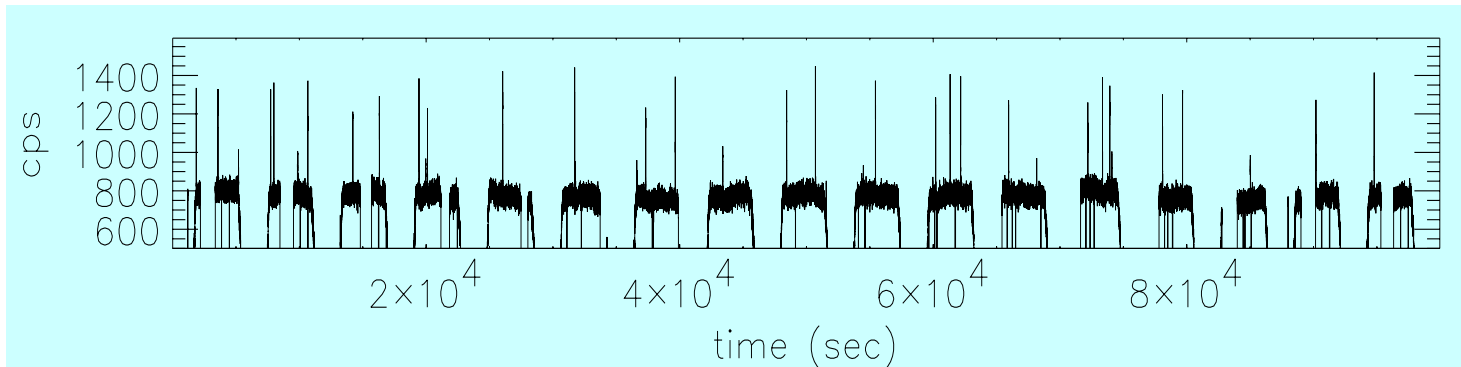
- Accretion at levels of 10^{-10} / $10^{-9} M_{\odot}/\text{yr}$ in CVs with P_{orb} between 1 and 10 hrs
- White dwarf is heated by compressional heating
- Envelope electron-degenerate \rightarrow pressure not a function of temperature
- Heating and pressure build-up ignites accreted hydrogen layer, first through pp, then through CNO cycle. T not high enough for helium burning (3α)
- Burning of freshly accreted material can continue for some time (~ 100 d) due to high T
- Nuclear network moderate: ~ 100 isotopes up to $A \approx 40$, $T_{\text{peak}} \approx 10^8$ K. All nuclear reaction rates based on experimental data. This is a significant contrast to X-ray bursts (500 isotopes, 1000s of reactions, incomplete experimental data)

Comparison Novae / X-ray bursts

	Novae	X-ray burst	Ratio
Galactic rate per year	40	10^4	10^3
R (km)	10^4	10	10^3
g_{surface} (cm s ⁻²)	10^{11}	10^{14}	10^3
Recurrence time	decades to megayears	hours to weeks	10^3
Duration	weeks to years	Minute to hours	10^3
Fluence (ergs)	10^{41-46}	10^{39-42}	10^4
Peak luminosity	Always Eddington	Up to Eddington	
T_{peak} (K) of burning	10^8	10^9	
kT_{eff} (keV)	0.05	2	$\sqrt{10^3}$
Isotopes	10^2 (A<40)	10^3 (A<100)	
Alpha	~0.01	~100	10^4
Mass burned (M_{\odot})	10^{-4}	10^{-12}	10^8
Density (g cm ⁻³)	10^4	10^6	10^2

X-ray bursts are far more frequent..

One day in the *Galactic bulge* → tens of X-ray bursts!



Critical envelope mass, recurrence time

$$P_{\text{crit}} \sim 10^{19} \text{Nm}^{-2}$$

$$P = \frac{F}{A} = \frac{GMM_{\text{env}}}{R^2} \frac{1}{4\pi R^2} \stackrel{R \propto M^{-1/3}}{\propto} \frac{MM_{\text{env}}}{R^4} \propto \frac{MM_{\text{env}}}{M^{-4/3}} = M^{7/3} M_{\text{env}}$$

$$\implies M_{\text{env}} \propto \frac{P_{\text{crit}}}{M^{7/3}} \quad M \uparrow \implies M_{\text{env}} \downarrow$$

$$M \sim 0.6M_{\odot} \implies M_{\text{crit}} = 5 \times 10^{-3} M_{\odot}$$

$$M \sim 1.3M_{\odot} \implies M_{\text{crit}} = 3 \times 10^{-5} M_{\odot}$$

$$\tau_{\text{rec}} = \frac{\langle M_{\text{crit}} \rangle}{\langle \dot{M} \rangle} \simeq 10^3 \dots 10^6 \text{y}$$

$\dot{M} \simeq 10^{-11} \dots 10^{-9} M_{\odot} \text{yr}^{-1}$

$$N_{\text{CN}} \simeq \frac{M_{2,i}}{\langle M_{\text{crit}} \rangle} \simeq 10^2 \dots 10^4$$

Stellar Transients /
Thermonuclear shell flashes II

Temperature vs accretion rate → Ignition mass vs accretion rate

- Heating due to compressional heating; source gravitational energy

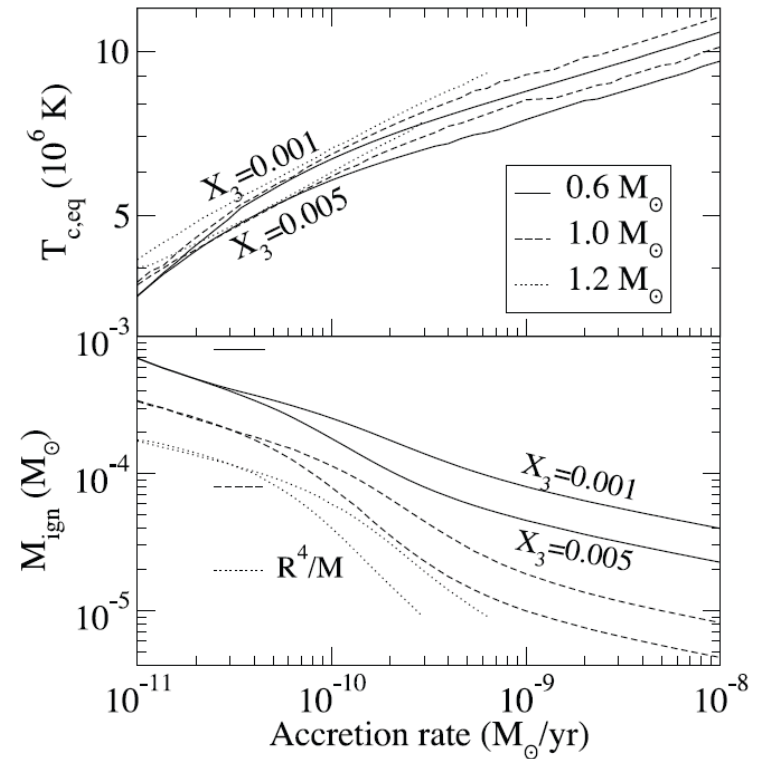
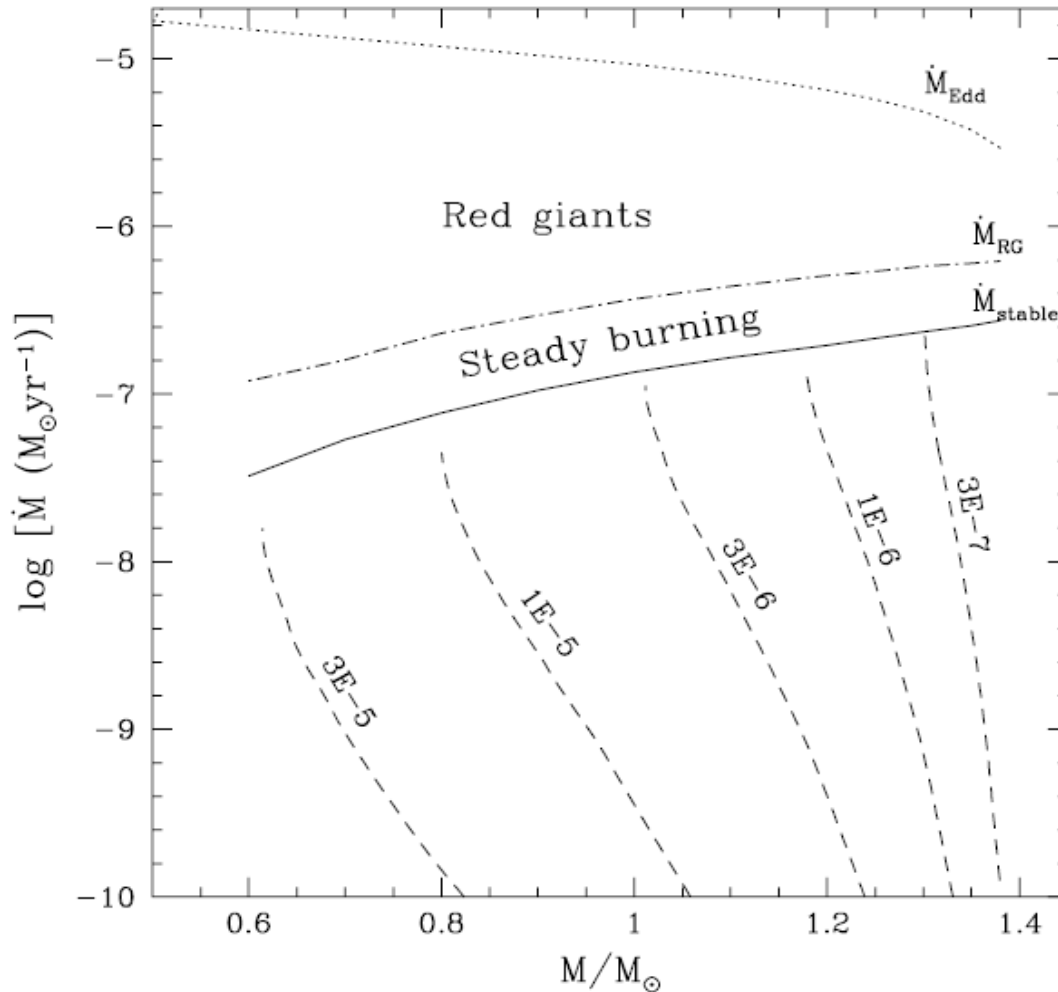


FIG. 8.—Equilibrium core temperature and mass of the accreted layer at CN ignition, M_{ign} , as a function of the time-averaged accretion rate $\langle \dot{M} \rangle$. Curves are shown for three WD masses M and two different mass fractions of ${}^3\text{He}$ at each value of M .

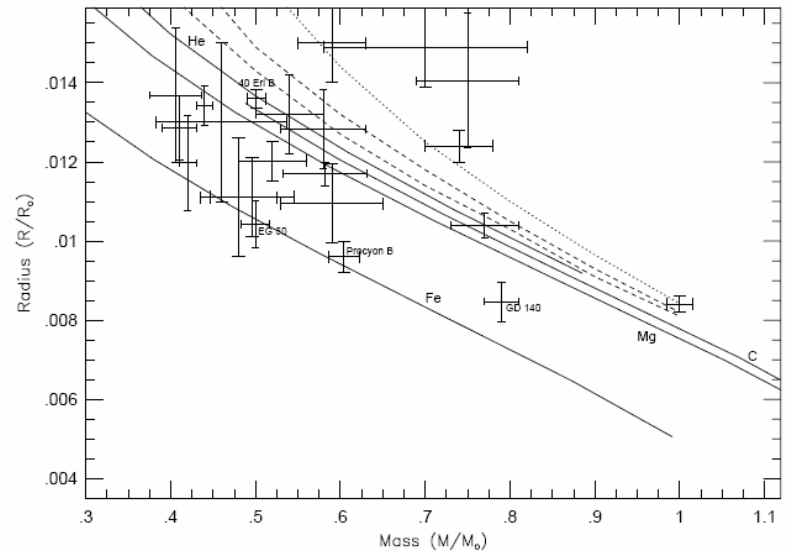
Stable and unstable burning on WD



Note that Eddington limit on WDs is 10^3 times higher than on NSs in terms of \dot{M} , but similar in terms of L . Reason: bigger WD \rightarrow only ~ 0.1 MeV/nucleon grav energy release in WD instead of 200 MeV/nucleon as in NS

Matter expulsion

- GM/R between 0.04 and 0.2 MeV/nucleon
- Nuclear energy 7 MeV/nucleon \rightarrow there is plenty of energy to expel each nucleon, but expulsion will only occur if the Eddington limit is reached \rightarrow Question: how fast does burning need to be before Eddington limit is reached?
- Eddington limit is $8.1 \times 10^{43} (M_{WD}/M_{\odot})$ MeV/s
- The Eddington limit can be reached if $1.2 \times 10^{43} (M_{WD}/M_{\odot})$ nucleons/s are burned and energy is radiated
- For a fuel mass of $10^{-5} M_{\odot} = 1.2 \times X \times 10^{52}$ H-nucleons the burning time should be less than 20 yrs to keep flux at Edd.
- Classical novae last shorter \rightarrow always super Eddington



Provencal et al. 2004

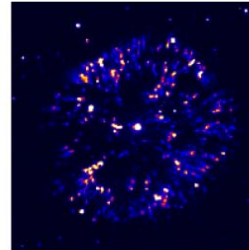
Nova remnants

- Most accreted mass expelled, after being enriched through mixing with deeper layers in WD
- Sometimes odd abundances in remnants \rightarrow excavation of CO or ONe WD proposed (meaning that WD mass decreases through classical novae!)
- A lot of H is found in ejecta (up to 50%) \rightarrow not all H is burnt
- ISM pollution: ~ 30 CNs/yr/Gal $\times 10$ Gyr $\times 10^{-5} M_{\odot} \approx 10^7 M_{\odot} \rightarrow$ not important, except:
- Some isotopes (^{13}C , ^{15}O , ^{17}Ne) 1000x overabundant \rightarrow may be important

Tim O'Brien's Gallery of Novae

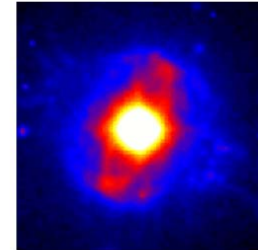
GK Persei

Outburst 1901



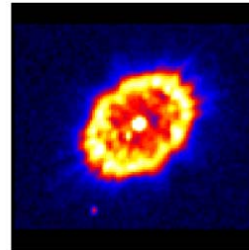
RR Pictoris

Outburst 1927



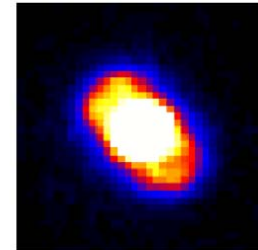
DQ Herculis

Outburst 1934



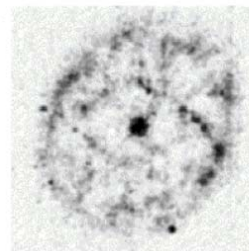
HR Delphini

Outburst 1967



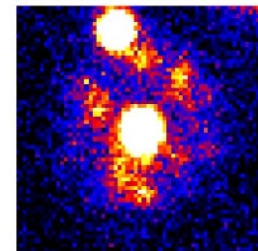
$\bar{\text{H}}$ Serpentis

Outburst 1970



V1500 Cyg

Outburst 1975



Light curve evolution

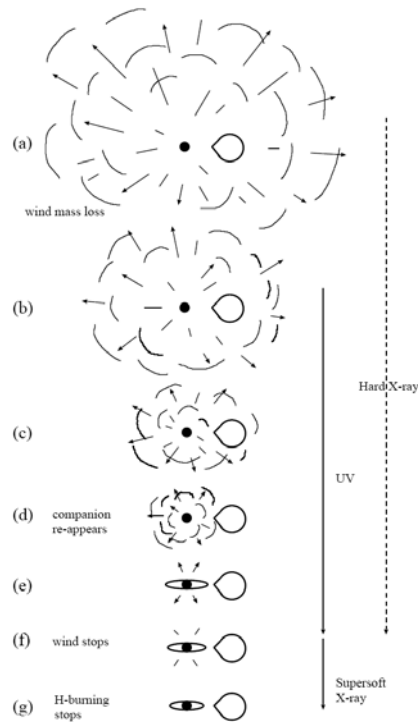


FIG. 1.— Evolution of nova outbursts: (a) after a nova explosion sets in, the photosphere expands greatly up to $\gtrsim 100 R_{\odot}$, and the companion star is engulfed deep inside the photosphere; (b) after the maximum expansion, the photospheric radius shrinks with time and free-free emission dominates the flux at relatively longer wavelengths; (c) a large part of the envelope matter is blown in the wind and the photosphere moves further inside; (d) the companion eventually emerges from the WD photosphere and an accretion disk may appear or reestablished again; (e) the photosphere further shrinks to a size of $\lesssim 0.1 R_{\odot}$; (f) the optically thick wind stops; (g) hydrogen nuclear burning stops and the nova enters a cooling phase. Hard X-rays may originate from internal shocks between ejecta (or a bow shock between ejecta and the companion) from stage (a) to (f) as indicated by a dashed line. The ultraviolet (UV) flux dominates from stage (b) to (f). Then the supersoft X-ray flux replaces the UV flux from stage (f) to (g).

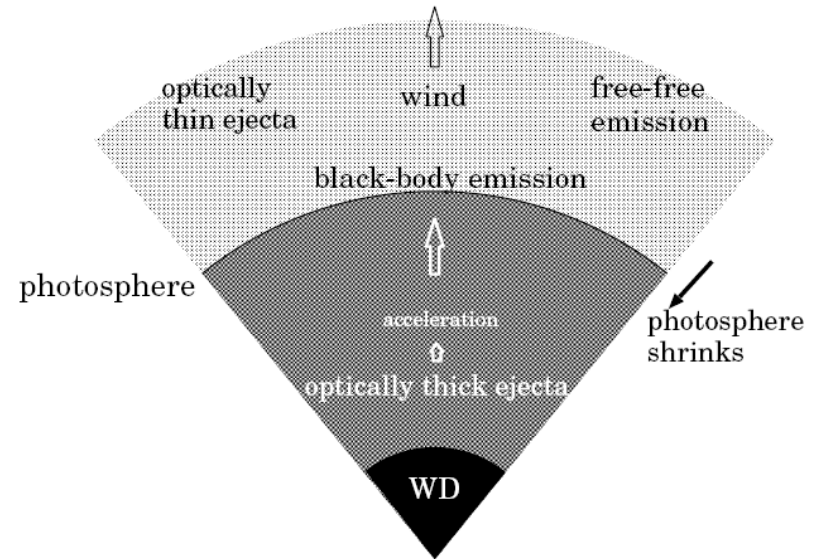
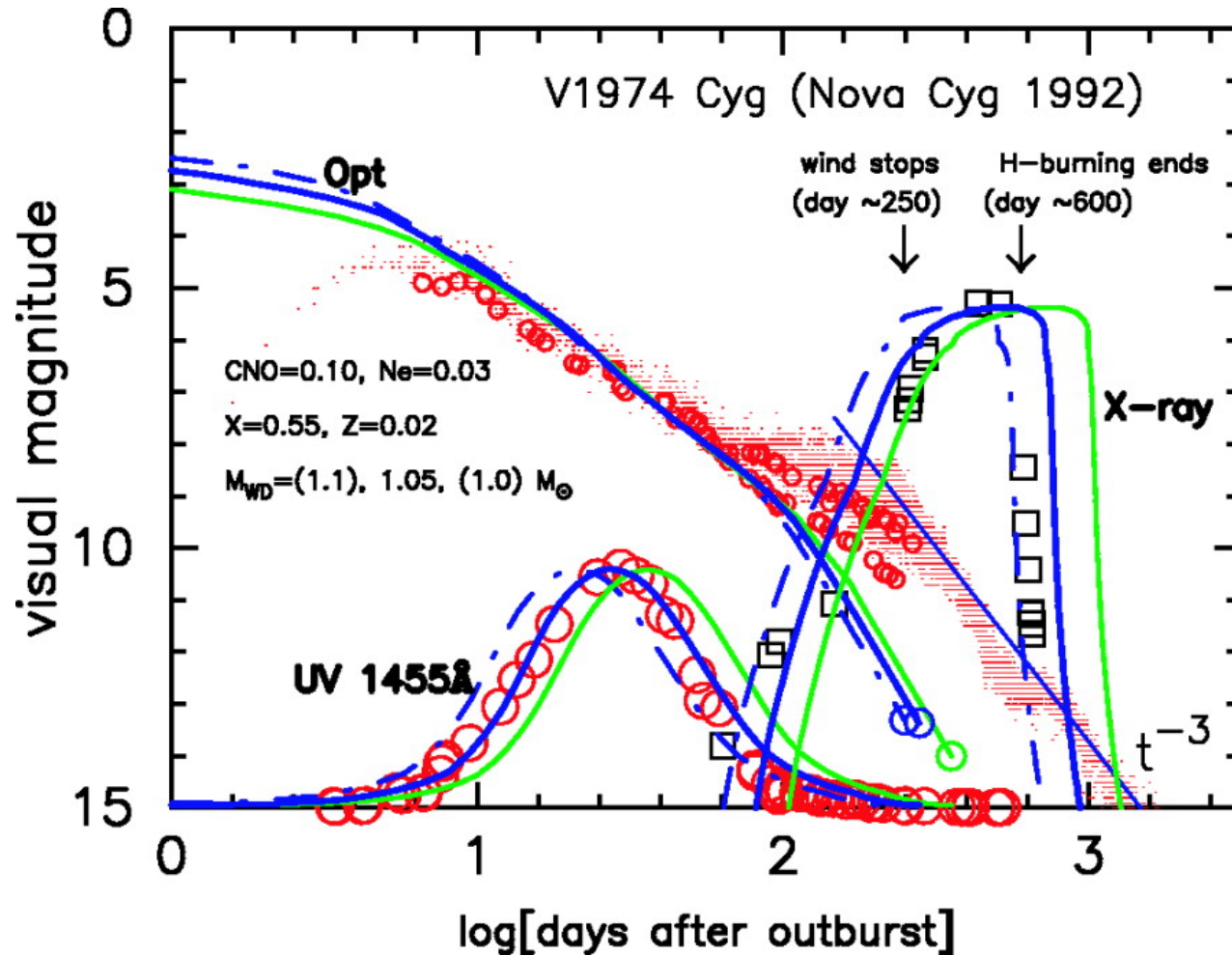


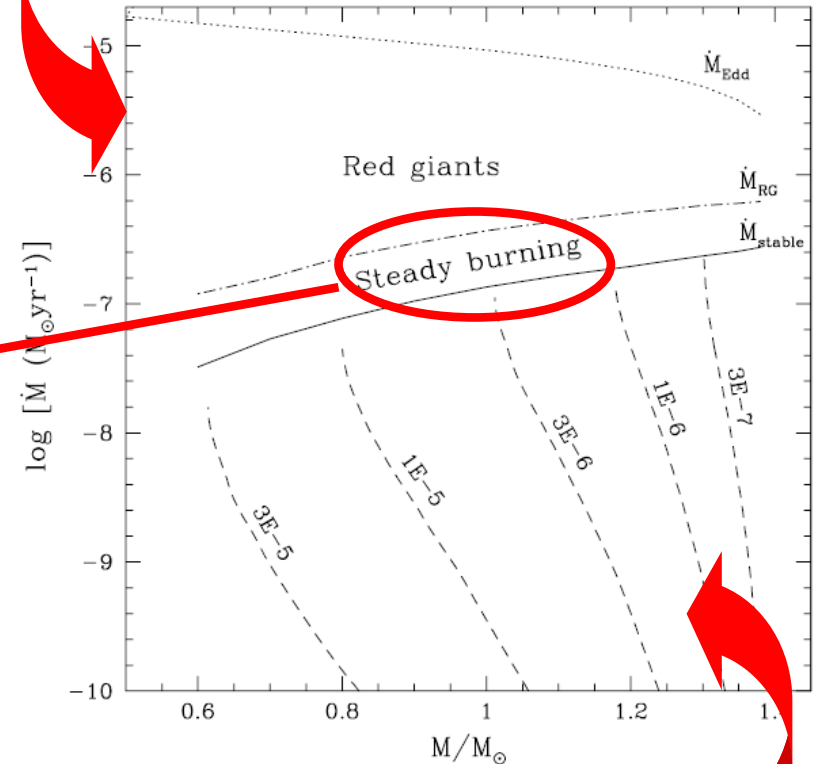
FIG. 2.— A schematic configuration of our nova ejection model: A large part of the initial envelope mass is ejected by the winds, which are accelerated deep inside the photosphere. After the optical maximum, that is, after the maximum expansion of the photosphere, the photosphere begins to shrink whereas the ejecta are expanding. The optically thin layer emits free-free radiation at relatively longer wavelengths while blackbody radiation from the photosphere dominates at shorter wavelengths.

Light curve @ 3 different wavelengths



Can we get a SN Ia?

- Can Chandrasekhar limit of $1.4 M_{\odot}$ be surpassed through accretion?
- What is needed in this 'single degenerate' channel?
 - WD
 - Heavy enough donor star, also leaving some room for non-conservative mass transfer
 - Transferred mass must be burned continuously instead of through novae, in order for it to be retained by the WD \rightarrow stringent limits on \dot{M} \rightarrow limits on donor (more massive, somewhat evolved)
- Candidate growing WDs: High \dot{M} (10^{38} erg/s), high T (tens of eVs) \rightarrow Super Soft Sources (SSS)

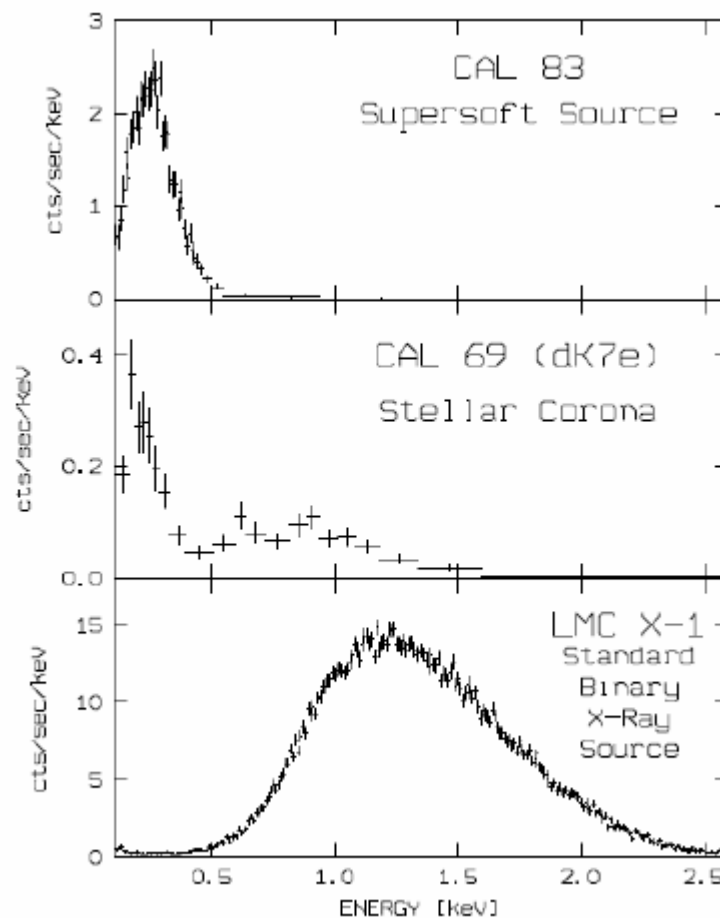


Accreted matter
lost in RG winds

Accreted matter
lost in novae

Super Soft Sources (10^{38} erg/s; tens of eVs)

- 18 in Galaxy and MCs
- 3 types
 - CVs with $M_{\text{donor}} > 1.5 M_{\odot}$
 - CVs with $M_{\text{donor}} < 1.5 M_{\odot}$
 - Symbiotics (supergiant+WD)
- 10^3 predicted, but many extinguished by interstellar dust \rightarrow look in other galaxies
- SSSs not seen in large numbers in external galaxies (~ 10 instead of 10^3)
- \rightarrow unlikely SN IA channel, at least if they are SSSs
- BUT: community not unanimous about this yet! To be continued..



Standard candle?

- There is a relationship between peak luminosity and nova duration: higher peak luminosity goes with faster novae. This can be understood through M_{WD} dependency
 - All novae are Eddington limited
 - $L_{\text{edd}} (\propto) M^{WD}$
 - $M_{\text{ign}} (\propto) M_{WD}^{-7/3}$
- Potential for use as standard candle and as distance indicator out to Virgo cluster (20 Mpc)

Thermonuclear shell flashes summary

- Occur on compact non-singular objects
- Need degenerate environment to keep pressure up to sustain nuclear burning
- Ignition mass dependent on \dot{M} , M and X
- Peak temperature dependent on X
- 4 important nuclear reaction chains: proton-proton chain (pp) , CNO cycle, triple alpha and rapid proton capture (rp). Latter two only in NSs
- Mass expulsion: from $<1\%$ (NS) to $\geq 100\%$ (WD)
- Light curve complications in WDs due to mass expulsion
- Stable hydrogen burning present in both NSs and WDs at similar accretion luminosities. The WD cases ('Super Soft Sources') are candidate SN Ia progenitors
- Agreement on general physics, NOT on details such as on accretion-induced Type Ia SNe.